

## ON THE SYNTHESIS OF NONUNIFORM PLANAR ANTENNA ARRAYS

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**Abstract :** *A synthesis of a nonuniform planar array is obtained by the method of orthonormalization of the base where the array factor  $f(\varphi)$  is referred. Some applications are carried out.*

### I. THEORETICAL SYNTHESIS

The problem of synthesizing a planar array can be solved like the one of linear arrays [1], [2].

The planar factor of an array consisting of  $N$  discrete elements is given by

$$f(\varphi) = \sum_{i=1}^N A_i \exp | jkr_i \cos(\varphi - \varphi_i) | \quad (1)$$

where  $k = 2\pi/\lambda$  and  $(r_i, \varphi_i)$  are the planar co-ordinates of the  $i$ th element. One can find the amplitudes  $A_i$  by orthonormalizing the base referring to  $f(\varphi)$ . Thus, the infinite set

$$\{\Phi_i(\varphi) = \exp | jkr_i \cos(\varphi - \varphi_i) |\}$$

will, on applying Schmidt's method, give the new set

$$\{\Psi_j(\varphi) = \sum_{i=1}^j C_j^{(i)} \exp | jkr_i \cos(\varphi - \varphi_i) |\} \quad (2)$$

the coefficients  $C_j^{(i)}$  of (2) will be determined with the aid of the following relations:

$$\Psi_1(\varphi) = \frac{\Phi_1(\varphi)}{|\langle \Phi_1, \Phi_2 \rangle|^{1/2}}$$

$$\Psi_n(\varphi) = \frac{\Phi_n(\varphi) - \sum_{j=1}^{n-1} \langle \Phi_1, \Psi_j, \rangle \Psi_j(\varphi)}{|\langle A_n, A_n \rangle|^{1/2}} \quad (3b)$$

where  $A_n$  is the numerator of (3b).

Taking into account that

$$\langle \Phi_n, \Psi_i \rangle = \int_0^{2\pi} \Phi_n(\varphi) \Psi_i^*(\varphi) d\varphi = 2\pi \sum_{i=1}^j C_i^{(j)} J_0(ke_{ni})$$

( $J_0$  is the zeroth Bessel function)

where

$$e_{ni} = \frac{r_n \sin \Phi_n - r_i \sin \Phi_i}{\sin \Phi_{ni}}$$

and

$$\cos \Phi_{ni} = \frac{r_n \cos \Phi_n - r_i \cos \Phi_i}{r_n \sin \Phi_n - r_i \sin \Phi_i}$$

considering (1) and (2) and solving (3a) and (3b), we get

$$C_n^{(n)} = \frac{1}{D_n}$$

$$C_k^{(n)} = -\frac{2\pi}{D_n} \sum_{j=1}^{n-1} C_k^{(j)} \left( \sum_{i=1}^j C_i^{(j)} J_0(ke_{ni}) \right)$$

$$D_n = \left| 2\pi - 4\pi^2 \sum_{j=1}^{n-1} \left( \sum_{i=1}^j C_i^{(j)} J_0(ke_{ni}) \right)^2 \right|^{1/2} \quad (4)$$

$$B_j = \int_0^{2\pi} f(\varphi) \Psi_j^*(\varphi) d\varphi$$

and

$$A_i = \sum_{j=i}^N B_j C_i^{(j)}$$

## II. MAXIMUM GAIN

To have a maximum gain in a certain direction  $\varphi_0$ , it is necessary for the radiation pattern to be an impulse function in this direction. Given that the array plane can be considered as the Oxy plane of an orthogonal system  $Oxyz$ , there follows that the maximum will appear in a direction  $\theta = \pi/2$ ,  $\varphi = \varphi_0$ .

The radiation pattern is

$$F(\varphi, \theta) = g(\varphi, \theta) \sum A_i e^{jkr_i \cos(\varphi - \varphi_i)} \quad (5)$$

On the array plane,

$$F(\varphi, \pi/2) = g(\varphi, \pi/2) \sum A_i \exp\{jkr_i \cos(\varphi - \varphi_i)\} \quad (6)$$

In agreement with the above discussion, it is necessary that

$$\left. \begin{aligned} g(\varphi_0, \pi/2) &= 1 \quad \text{and} \\ F(\varphi, \pi/2) &= \delta(\varphi - \varphi_0) \end{aligned} \right\} \quad (7)$$

Then

$$\begin{aligned} B_j &= \int_0^{2\pi} f(\varphi) \Psi_j^*(\varphi) d\varphi = \int_0^{2\pi} \frac{F(\varphi, \pi/2)}{g(\varphi, \pi/2)} \sum \Psi_i^*(\varphi) d\varphi \\ &= \int_0^{2\pi} \frac{\delta(\varphi - \varphi_0)}{g(\varphi, \pi/2)} \Psi_i^*(\varphi) d\varphi = \Psi_i^*(\varphi_0) \end{aligned} \quad (8)$$

The gain will be

i) for isotropic elements ( $g(\varphi, \theta) = 1$ )

$$\begin{aligned} G &= \frac{|F(\varphi_0, \pi/2)|^2}{1/4\pi \int_0^\pi \int_0^{2\pi} |F(\varphi, \theta)|^2 \sin\theta d\varphi d\theta} = \frac{|f(\varphi_0)|^2}{1/4\pi \int_0^\pi \int_0^{2\pi} |f(\varphi)|^2 \sin\theta d\varphi d\theta} \\ &= 2\pi \frac{\sum B_i B_i^{*2}}{\sum B_i B_i^*} = 2\pi \sum B_i B_i^* \end{aligned} \quad (9)$$

ii) For parallel dipoles, normal to the array plane,

$$(g(\varphi, \theta) = \sqrt{1 - \sin^2 \theta \cos^2 \varphi}),$$

$$G = \frac{|F(\varphi_0, \pi/2)|^2}{\int_0^\pi \int_0^{2\pi} |F(\varphi, \theta)|^3 \sin \theta d\varphi d\theta} =$$

$$= 2\pi \frac{|\sum B_i B_i^*|^2}{\sum B_i B_i^* - 2/3 \sum \sum A_i A_j \{J_0(ke_{ij}) - \cos \varphi_{ij} J_1(ke_{ij})\}} \quad (10)$$

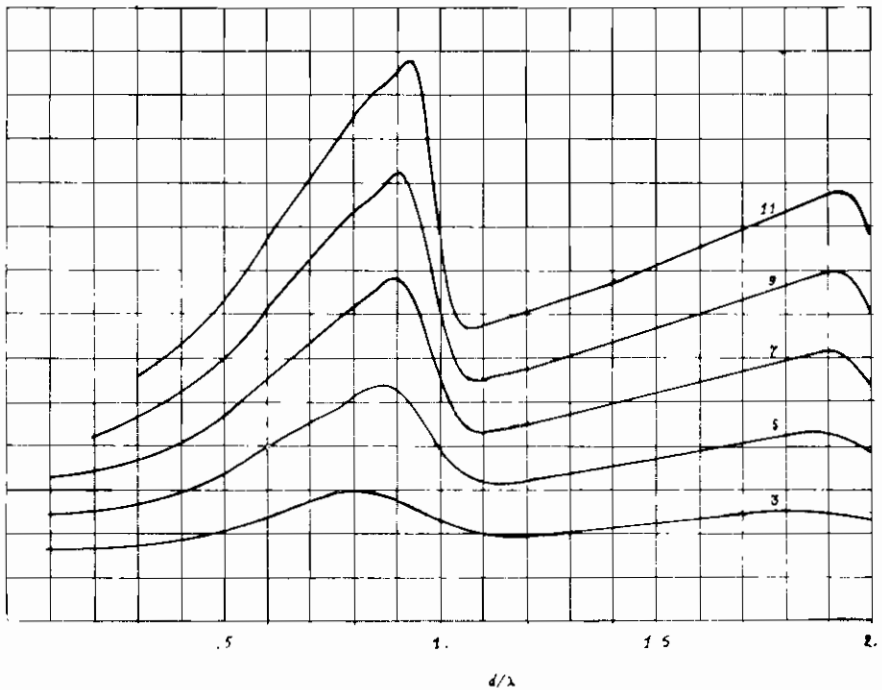


Fig. 1  
Optimum gain of axial arrays of crossed dipoles  
( $N = 3, 5, 7, 9, 11$ )

Figs (1) and (2) represent the gain for two cases: The first regards a uniform circular array of isotropic elements, and the second a uniform linear array of dipoles; for both cases, one cannot fail to remark the coincidence with the nomographs given by Cheng [4], Tai [5] respectively.

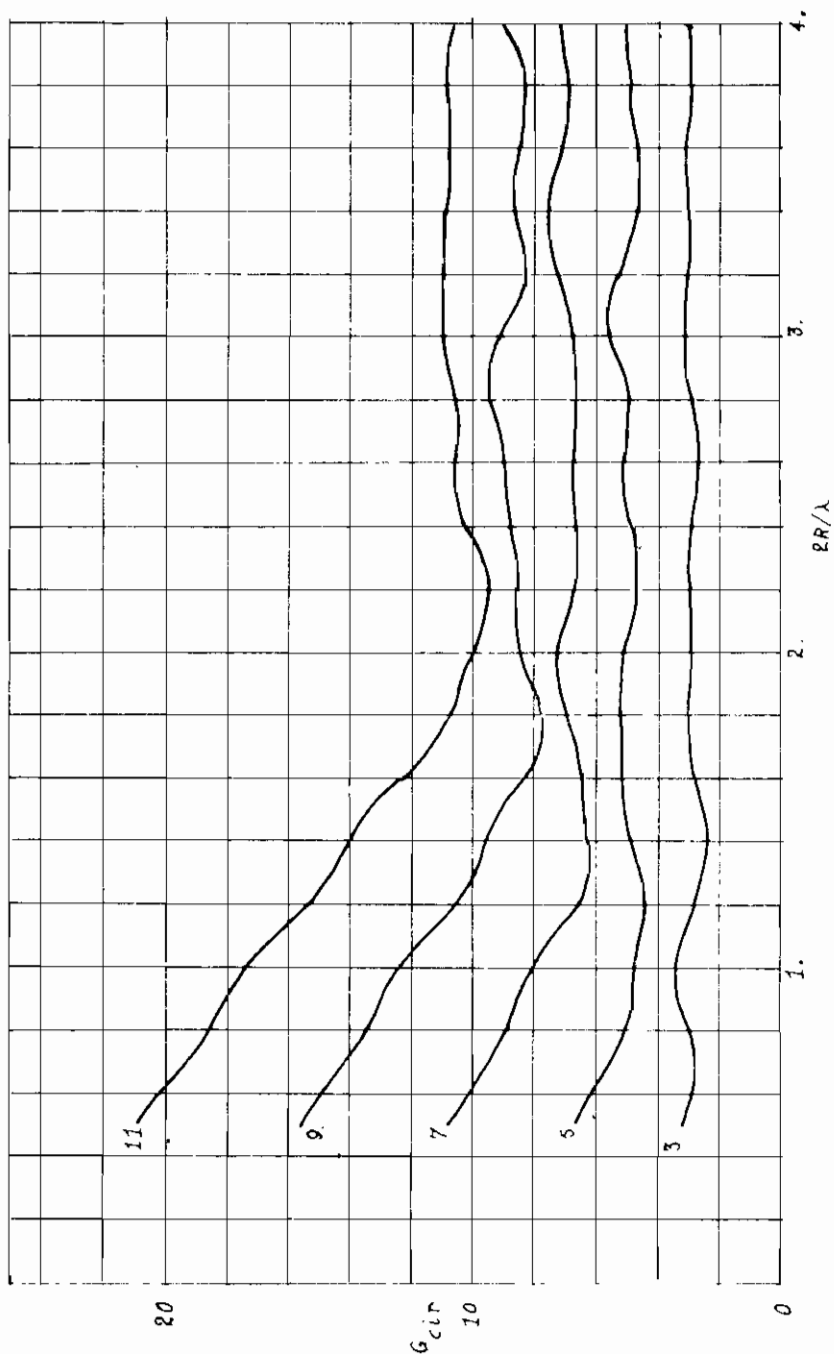


Fig. 2 Optimum gain of circular arrays of isotropic elements ( $N = 3, 5, 7, 9, 11$ )

### III. A DIRECTIVE RADIATION PATTERN

A directive radiation pattern can be expressed as

$$f(\varphi) = \exp\{p\cos(\varphi - a)\}\cos\{q\sin(\varphi - a)\} \quad (11)$$

Maximum occurs at  $\varphi = a$ , while the first null when

$$q\sin(\varphi_0 - a) = \pi/2 \Rightarrow \varphi = \frac{\pi}{2\sin(\varphi_0 - a)} \quad (12)$$

A secondary maximum can be defined by zeroing the derivative of  $f(\varphi)$ , that is

$$P\sin(\varphi_1 - a)\cos|q\sin(\varphi_1 - a)| + q\cos(\varphi_1 - a)\sin|q\sin(\varphi_1 - a)| = 0$$

The side lobe level of the radiation pattern is

$$M = \frac{F(a)}{F(\varphi)} = \frac{\exp\{P|1 - \cos(\varphi_1 - a)|\}}{\cos|q\sin(\varphi_1 - a)|} \quad (14)$$

For  $a = 0$ , one gets

$$\begin{aligned} q &= \frac{\pi}{2\sin\varphi_0} \\ P &= -\frac{\pi}{2\sin\varphi_0} \cdot \frac{\tan\left|\frac{\pi\sin\varphi_1}{2\sin\varphi_0}\right|}{\tan\varphi_1} \\ M &= \frac{\exp\{P(1 - \cos\varphi_1)\}}{\cos(q\sin\varphi_1)} \end{aligned} \quad (15)$$

With the aid of (15) we can get  $P$ ,  $q$  and  $M$  versus  $\varphi_1$ ,  $\varphi_0$ , as shown in figs (3a,b,c).

From (11) we can get the values of  $B_j$  as follows:

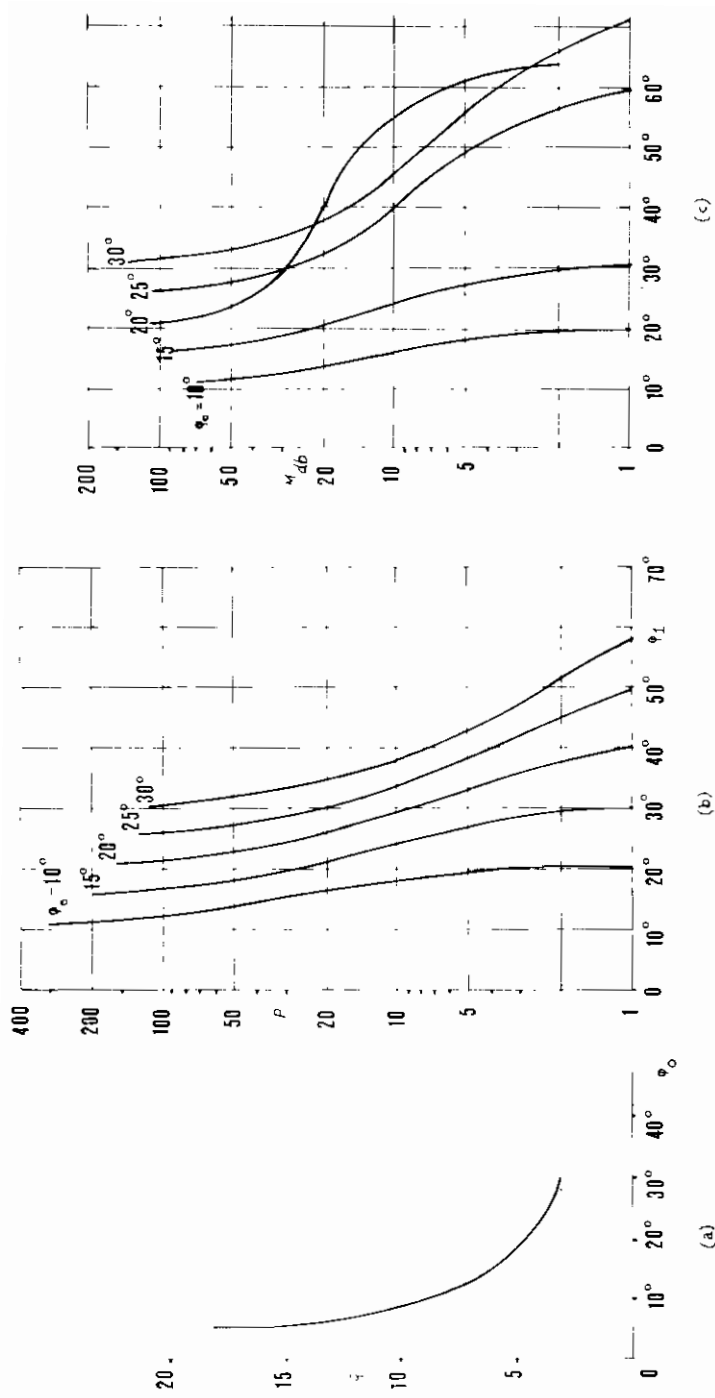


Fig. 3 g,  $P$  and  $M_{1db}$  versus  $\phi_0, \phi_1$

$$B_i = \pi \sum_{l=1}^j C_l^{*(Q)} | I_o (\sqrt{b_l - jC_l}) + I_o (\sqrt{d_l - jf_l}) | \quad (16)$$

where

$$\begin{aligned} b_i &= P^2 - g_i^2 & , & & C_i &= -2Pg_i \cos(\Phi_i - \Psi_i) \\ d_i &= P^2 - h_i^2 & & & f_i &= -2Ph_i \cos(\Phi_i - Z_i) \end{aligned}$$

$$g_i = \frac{q \cos \alpha + kr_i \sin \Phi_i}{\sin \Psi_i} \quad , \quad h_i = \frac{q \cos \alpha - kr_i \sin \Phi_i}{\sin Z_i}$$

$$\cos \Psi_i = \frac{q \sin \alpha + kr_i \cos \Phi_i}{q \cos \alpha + kr_i \sin \Phi_i} \quad , \quad \cos Z_i = \frac{q \sin \alpha - kr_i \cos \Phi_i}{q \cos \alpha - kr_i \sin \Phi_i}$$

#### IV. CONCLUSION

As shown by the preceding discussion, it is possible to synthesize non-uniform planar arrays by the orthonormalization method. The maximum gain can be easily found on replacing the radiation pattern by the  $\delta$  — Dirac function.

By choosing a directive radiation pattern it is possible to synthesize an array having pre-defined i) angle of maximum, ii) width between the first nulls and iii) Side lobe level.

All these apply equally well to linear arrays, too.



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ΠΕΡΙΛΗΨΙΣ

ΠΕΡΙ ΤΗΣ ΣΥΝΘΕΣΕΩΣ ΜΗ ΟΜΟΓΕΝΩΝ ΕΠΙΠΕΔΩΝ  
ΣΤΟΙΧΕΙΟΚΕΡΑΙΩΝ

ὑπὸ

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